## Complete control over reflected fields with gap surface plasmons Sergey I. Bozhevolnyi Centre for Nano Optics, University of Southern Denmark

#### Introduction:

- Gap surface plasmons (GSPs)
- GSP resonators and their arrays
- Black and coloured metasurfaces
- Phase-gradient metasurfaces
- Independent polarization control
- Analog computing
- Polarimetry by metasurfaces
- Stealth technologies
- Instead of outlook





Sergey I. Bozhevolnyi

5 µm

## Intro: gap surface plasmons

The GSP mode **confinement** and propagation constant **increase** with the decrease in the gap width (no cutoff!), storing field in metal (the propagation loss increases).



# Confined GSP modes should be efficiently reflected by the terminations forming nano-resonators!

#### **General properties of slow-plasmon resonant** nanostructures: nano-antennas and resonators

20 August 2007 / Vol. 15, No. 17 / OPTICS EXPRESS 10869 Sergey I. Bozhevolnyi and Thomas Søndergaard

Department of Physics and Nanotechnology, Aalborg University, Skjernvej 4A, DK-9220 Aalborg Øst, Denmark



magnetic dipole:

P@N KCL seminar (London, October 5, 2016)







PHYSICAL REVIEW B 79, 035401 (2009)

#### Gap plasmon-polariton nanoresonators: Scattering enhancement and launching of surface plasmon polaritons

Jesper Jung\* and Thomas Søndergaard Department of Physics and Nanotechnology, Aalborg University, Skjernvej 4A, DK-9220 Aalborg Øst, Denmark

Sergey I. Bozhevolnyi Institute of Sensors, Signals and Electrotechnics (SENSE), University of Southern Denmark, Niels Bohrs Allé 1, DK-5230 Odense M, Denmark





Fig. 1. (a) Sketch of a two-dimensional gold-SiO<sub>2</sub>-gold resonator surrounded by air. The incident field is TM-polarized and propagates along the y-axis. (b) Scattering and absorption cross sections (CS) normalized to the width w of the gold strips for three different combinations of w and d. The strip thickness is fixed at t = 30 nm.

# GSP resonances can be strongly absorbing or strongly scattering, depending on the gap width!

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Fig. 2. (a,b) Contribution to the scattering cross section (SCS) from the lowest order multipoles (ED=electric dipole, MD=magnetic dipole, EQ=electric quadrupole). SCS is normalized with the resonator width w, and t = 30 nm [see Fig. 1(a)]. Note that in order to compare the relative contributions to the scattering from MD and EQ, we choose the center of mass as the coordinate origin. (c,d) Electric field enhancement at the ED mode ( $\lambda = 585$  nm) and GSP mode ( $\lambda = 800$  nm). The color bars are chosen as to emphasize the mode profiles rather than the high electrostatic field enhancement at the corners. Arrows indicate the

# The main contribution to GSP resonances comes from the magnetic dipole resonances!

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Opt. Express 21, 27438 (2013). Syddansk Universitet Sergey



# GSP resonators can be designed with continuous bottom metal and dielectric spacer layers!

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Opt. Express 19, 19310 (2011).

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## **Plasmonic black metasurfaces**

## Efficient absorption of visible radiation by gap plasmon resonators

Michael G. Nielsen,<sup>1,\*</sup> Anders Pors,<sup>2</sup> Ole Albrektsen,<sup>1</sup> and Sergey I. Bozhevolnyi<sup>1</sup>



Fig. 1. (a) Schematic representation of the quadratic unit-cell of size A, and containing four gold nanoparticles with different diameters d. (b) Schematic representation of the layered sample structure. (c) SEM-image of the fabricated CL-GPRs with particle diameters  $d_1$  =



Fig. 3. (a) Schematic representation of the unit-cell comprising four gold nanoparticles with different diameters and nine additional small gold nanoparticles with identical diameter. The unit-cell size is  $\Lambda$  = 340nm. (b) SEM-image of the fabricated CL-GPRs with particle diameters

#### Average absorption: ~ 94% (400 - 750 nm) and ~ 89% (400 - 850 nm)

Reflection

film coated with ~20nm SiO2.

coated with ~20nm SiO2. (d) Reflection spectra of a similar CL-GPR array, but with unit-cell

#### **GSP** resonances ensure very efficient absorption!

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4 June 2012 / Vol. 20, No. 12 / OPTICS EXPRESS 13311

## **Plasmonic multicolored metasurfaces**



# GSP resonances, when properly designed, can be used for sub-diffraction colour printing!

European Research Council Established by the European Commission SYDDANSK UNIVERSIT

## GSP resonator arrays: from amplitude to phase control

y j<sup>Z</sup> x t t t s

Opt. Express **21**, 2942 (2013)

 $t = 50 \text{ nm}, L_x = 138 \text{ nm}, L_y = 105 \text{ nm}$  $E_{in} = (E_x, 0, 0)$ 

![](_page_8_Figure_4.jpeg)

## **GSP** resonances, when properly designed, can ensure either efficient absorption or reflection!

European Research Council Established by the European Commission

## GSP resonator arrays: from amplitude to phase control

![](_page_9_Figure_1.jpeg)

# GSP resonances, when carefully designed, result also in strong (reflected light) phase variation!

## GSP resonators: from amplitude to phase control

![](_page_10_Figure_1.jpeg)

## Extra large phase variation is due to drastic changes in path configurations of reflected light!

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## **GSP-based perfectly reflecting half-plates**

![](_page_11_Figure_1.jpeg)

## GSP resonator parameters control the phase difference between different polarizations!

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## **Phase-gradient metasurfaces**

#### Historical perspective: from blazed gratings to phase-gradient surfaces

![](_page_12_Figure_2.jpeg)

## **Prehistoric phase-gradient metasurfaces**

1424 OPTICS LETTERS / Vol. 26, No. 18 / September 15, 2001

#### Pancharatnam–Berry phase in space-variant polarization-state manipulations with subwavelength gratings

Ze'ev Bomzon, Vladimir Kleiner, and Erez Hasman

Optical Engineering Laboratory, Faculty of Mechanical Engineering, Technion—Israel Institute of Technology, Haifa 32000, Israel

![](_page_13_Figure_5.jpeg)

Fig. 3. (a) Magnified geometry of the grating for converting circular polarization into azimuthal polarization and (b) experimental measurement of the local azimuthal angle  $\psi$ .

![](_page_13_Figure_7.jpeg)

Fig. 4. Measured and calculated cross sections of the far-field images for (a) the in-phase and (b) the antiphase azimuthal polarization. Inset, experimental intensity distributions.

We realized a Lee-type binary grating ... The metal stripes consisted of 10 of nm Ti and 60 of nm Au, which were deposited onto 500-mm-thick GaAs wafers by photolithography and a lift-off technique. ..

We illuminated the grating with circularly polarized light at a wavelength of 10.6  $\mu$ m.

uropean Research Council labilithed by the European Commission

![](_page_13_Picture_12.jpeg)

## Prehistoric phase-gradient metasurfaces

Metal metasurfaces:

1. "*Pancharatnam-Berry phase in space-variant polarization state manipulations with subwavelength gratings*", Z. Bomzon, V. Kleiner, E. Hasman, Opt. Lett. **26**, 1424-1426 (2001).

2. "Formation of radially and azimuthally polarized light using space-variant subwavelength metal stripe gratings", Z. Bomzon, V. Kleiner, E. Hasman, Appl. Phys. Lett. **79**, 1587-1589 (2001).

Dielectric metasurfaces:

3. "Space-variant Pancharatnam-Berry phase optical elements with computer-generated subwavelength gratings", Z. Bomzon, G. Biener, V. Kleiner, E. Hasman, Opt. Lett. 27, 1141-1143 (2002).

4. "*Formation of helical beams by use of Pancharatnam-Berry phase optical elements*", G. Biener, A. Niv, V. Kleiner, and E. Hasman, Opt. Lett. **27**, 1875-1877 (2002).

5. "Polarization dependent focusing lens by use of quantized Pancharatnam-Berry phase diffractive optics". E. Hasman, V. Kleiner, G. Biener, and A. Niv, Appl. Phys. Lett. 82, 328-330.

### There is nothing new under the sun – from Ecclesiastes 1:9

## Why phase-gradient metasurfaces?

Historical perspective: from blazed gratings to phase-gradient surfaces

![](_page_15_Figure_2.jpeg)

whereas thickness of metasurfaces is NOT!

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![](_page_15_Picture_5.jpeg)

![](_page_15_Picture_6.jpeg)

## **GSP-based phase-gradient metasurfaces**

![](_page_16_Figure_1.jpeg)

15°

## **GSP-based phase-gradient metasurfaces**

![](_page_17_Figure_1.jpeg)

## **GSP-based phase-gradient metasurfaces**

![](_page_18_Figure_1.jpeg)

# The crucial features of GSP metasurfaces:

![](_page_19_Figure_1.jpeg)

**Reflected amplitude and phase for x-polarization** 

#### Two polarizations can be controlled independently!

![](_page_19_Picture_5.jpeg)

![](_page_20_Figure_0.jpeg)

## GSP-based phase-gradient metasurfaces: <u>Independent control of two polarizations!</u>

![](_page_21_Figure_1.jpeg)

![](_page_21_Figure_2.jpeg)

Metasurface functionality is robust against fabrication tolerances.

Lower efficiency compared to theory.

Sci. Rep. **3**, 2155 (2013).

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## GSP-based phase-gradient metasurfaces: <u>Independent control of two polarizations!</u>

![](_page_22_Figure_1.jpeg)

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![](_page_22_Figure_2.jpeg)

#### Optical characterization:

![](_page_22_Figure_4.jpeg)

(A) Diffraction spots. (B and C) Diffraction efficiencies. Curves: simulations; Markers: Experiment.

Good agreement between measurements and simulations (include additional losses).

Sci. Rep. **3**, 2155 (2013).

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## GSP-based phase-gradient metasurfaces: <u>Independent</u> control of two polarizations!

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![](_page_23_Figure_1.jpeg)

![](_page_23_Figure_2.jpeg)

(A) Diffraction spots. (B and C) Diffraction efficiencies.
 Curves: simulations; Markers: Experiment.

Functionality is robust against fabrication imperfections.

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Sci. Rep. **3**, 2155 (2013).

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# GSP-based phase-gradient metasurfaces: what can be next?

Case 1

Case 2

Case 3

![](_page_24_Figure_4.jpeg)

![](_page_24_Figure_5.jpeg)

### Phase-gradient metasurfaces: <u>Unidirectional & polarization-controlled</u> SPP couplers!

#### Polarization-Controlled Tunable Directional Coupling of Surface Plasmon Polaritons

Jiao Lin,<sup>1,2</sup>\* J. P. Balthasar Mueller,<sup>1</sup>\* Qian Wang,<sup>3</sup> Guanghui Yuan,<sup>3</sup> Nicholas Antoniou,<sup>4</sup> Xiao-Cong Yuan,<sup>5</sup> Federico Capasso<sup>1</sup>†

![](_page_25_Figure_3.jpeg)

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SCIENCE VOL 340 19 APRIL 2013

![](_page_25_Figure_6.jpeg)

Lingling Huang\*<sup>1,2</sup>, Xianzhong Chen\*<sup>1</sup>, Benfeng Bai<sup>2</sup>, Qiaofeng Tan<sup>2</sup>, Guofan Jin<sup>2</sup>, Thomas Zentgraf<sup>3</sup> and Shuang Zhang<sup>1</sup>

![](_page_25_Picture_8.jpeg)

### GSP-based phase-gradient metasurfaces: <u>Unidirectional & polarization-controlled</u> SPP couplers!

1D configuration:

![](_page_26_Figure_2.jpeg)

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### **GSP-based phase-gradient metasurfaces: Unidirectional & polarization-controlled SPP couplers!**

2D configuration:

![](_page_27_Figure_2.jpeg)

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GSP-based phase-gradient metasurfaces: <u>Unidirectional & polarization-controlled</u> SPP couplers! Experimental verification:

![](_page_28_Figure_1.jpeg)

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Light: Science & Applications, 2014

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GSP-based phase-gradient metasurfaces: <u>Unidirectional & polarization-controlled</u> SPP couplers! Experimental verification:

![](_page_29_Figure_1.jpeg)

erc

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Light: Science & Applications, 2014

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# GSP-based phase-gradient metasurfaces: what else can be done?

Case 1

Case 2

Case 3

![](_page_30_Picture_4.jpeg)

GSP-based gradient metasurfaces can efficiently and independently control orthogonal polarizations, including coupling to surface excitations!

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## Metasurfaces: performing mathematical operations!

![](_page_31_Figure_1.jpeg)

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![](_page_31_Picture_3.jpeg)

![](_page_31_Picture_4.jpeg)

# Optical holography with phase- and amplitude-gradient metasurfaces

![](_page_32_Figure_1.jpeg)

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![](_page_32_Picture_3.jpeg)

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![](_page_32_Picture_4.jpeg)

CCD Y-Polarizer  $Z = 0 \mu m$   $Z = 10 \mu m$   $Z = 15 \mu m$   $Z = 15 \mu m$  X Ar/Kr laserX-Polarizer

![](_page_33_Figure_1.jpeg)

#### <u>Main challenge:</u> obtain the opposite (constant) phase for the same (spatially varying) reflection amplitude!

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![](_page_33_Picture_4.jpeg)

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x (µm)

![](_page_34_Figure_1.jpeg)

## **GSP-based metasurfaces can do the job!**

## In practice, it is going to be very difficult!

![](_page_34_Picture_5.jpeg)

Nano Letters 15, 791 (2015)

**Differentiator configuration:** 

Integrator configuration:

![](_page_35_Figure_3.jpeg)

![](_page_35_Picture_4.jpeg)

![](_page_36_Figure_1.jpeg)

![](_page_37_Figure_1.jpeg)

### **GSP-based metasurfaces can do the job!**

![](_page_37_Picture_4.jpeg)

![](_page_37_Picture_5.jpeg)

## GSP-based phase-gradient metasurfaces for instantaneous SOP analysis

Monochromatic plane wave propagating along the *z*-axis:

$$\vec{E}(\vec{r}) = \begin{pmatrix} A_x \\ A_y e^{i\delta} \end{pmatrix} \exp(ikz); \ A_x > 0, \ A_y > 0$$

State of polarization (SOP) is completely determined by  $A_x/A_y$  and  $\delta$ : only two quantities have to be measured!

Main challenge: realization of simultaneous one-step measurements of all polarization contrasts!

$$s_{1} = A_{x}^{2} - A_{y}^{2},$$

$$s_{2} = 2A_{x}A_{y}\cos\delta = A_{a}^{2} - A_{b}^{2}, \quad (\hat{\mathbf{a}}, \hat{\mathbf{b}}) = \frac{1}{\sqrt{2}}(\hat{\mathbf{x}} + \hat{\mathbf{y}}, -\hat{\mathbf{x}} + \hat{\mathbf{y}})$$

$$s_{3} = 2A_{x}A_{y}\sin\delta = A_{r}^{2} - A_{l}^{2}, \quad (\hat{\mathbf{r}}, \hat{\mathbf{l}}) = \frac{1}{\sqrt{2}}(\hat{\mathbf{x}} + i\hat{\mathbf{y}}, \hat{\mathbf{x}} - i\hat{\mathbf{y}})$$

$$(\hat{\mathbf{r}}, \hat{\mathbf{l}}) = \frac{1}{\sqrt{2}}(\hat{\mathbf{x}} + i\hat{\mathbf{y}}, \hat{\mathbf{x}} - i\hat{\mathbf{y}})$$

$$(\hat{\mathbf{r}}, \hat{\mathbf{l}}) = \frac{1}{\sqrt{2}}(\hat{\mathbf{x}} + i\hat{\mathbf{y}}, \hat{\mathbf{x}} - i\hat{\mathbf{y}})$$

$$(\hat{\mathbf{r}}, \hat{\mathbf{l}}) = \frac{1}{\sqrt{2}}(\hat{\mathbf{x}} + i\hat{\mathbf{y}}, \hat{\mathbf{x}} - i\hat{\mathbf{y}})$$

$$(\hat{\mathbf{r}}, \hat{\mathbf{l}}) = \frac{1}{\sqrt{2}}(\hat{\mathbf{x}} + i\hat{\mathbf{y}}, \hat{\mathbf{x}} - i\hat{\mathbf{y}})$$

$$(\hat{\mathbf{r}}, \hat{\mathbf{l}}) = \frac{1}{\sqrt{2}}(\hat{\mathbf{x}} + i\hat{\mathbf{y}}, \hat{\mathbf{x}} - i\hat{\mathbf{y}})$$

$$(\hat{\mathbf{r}}, \hat{\mathbf{l}}) = \frac{1}{\sqrt{2}}(\hat{\mathbf{x}} + i\hat{\mathbf{y}}, \hat{\mathbf{x}} - i\hat{\mathbf{y}})$$

$$(\hat{\mathbf{r}}, \hat{\mathbf{l}}) = \frac{1}{\sqrt{2}}(\hat{\mathbf{x}} + i\hat{\mathbf{y}}, \hat{\mathbf{x}} - i\hat{\mathbf{y}})$$

$$(\hat{\mathbf{r}}, \hat{\mathbf{l}}) = \frac{1}{\sqrt{2}}(\hat{\mathbf{x}} + i\hat{\mathbf{y}}, \hat{\mathbf{x}} - i\hat{\mathbf{y}})$$

$$(\hat{\mathbf{r}}, \hat{\mathbf{l}}) = \frac{1}{\sqrt{2}}(\hat{\mathbf{x}} + i\hat{\mathbf{y}}, \hat{\mathbf{x}} - i\hat{\mathbf{y}})$$

$$(\hat{\mathbf{r}}, \hat{\mathbf{r}}) = \frac{1}{\sqrt{2}}(\hat{\mathbf{r}} + i\hat{\mathbf{y}}, \hat{\mathbf{r}} - i\hat{\mathbf{y}})$$

$$(\hat{\mathbf{r}}, \hat{\mathbf{r}}) = \frac{1}{\sqrt{2}}(\hat{\mathbf{r}} + i\hat{\mathbf{y}}, \hat{\mathbf{r}} - i\hat{\mathbf{y}})$$

Optica 2 P@N KCL

# GSP-based all-polarization sensitive birefringent metasurfaces

**Targeted operation principle:** 

![](_page_39_Picture_2.jpeg)

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![](_page_39_Picture_4.jpeg)

# GSP-based all-polarization sensitive birefringent metasurfaces

#### **Design considerations:**

![](_page_40_Figure_2.jpeg)

# GSP-based all-polarization sensitive birefringent metasurfaces

#### **Fabrication and characterization:**

![](_page_41_Figure_2.jpeg)

## GSP-based all-polarization sensitive birefringent metasurfaces

#### **Characterization results:**

![](_page_42_Figure_2.jpeg)

Diffraction contrast deviations of up to 0.15 are ascribed to fabrication imperfections, and to imperfections in optical components (in particular, in the quarter-wave plate).

Optica **2**, 716 (2015) P@N KCL seminar (London, October 5, 2016)

![](_page_42_Figure_5.jpeg)

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#### **All-polarization birefringent metagrating** $|x\rangle$ 10 **Characterization results:** |y > 10 $|a\rangle$ 0 $|b\rangle$ 0 $\lambda = 800 \text{ nm}$ $|r\rangle$ 0 $|1\rangle$ 01 news & views NATURE PHOTONICS | VOL 9 | NOVEMBER 2015 |

#### METASURFACES

## Simultaneous Stokes parameters

Techniques for determining Stokes parameters, which fully define the polarization state of a wave, require multiplemeasurements, thus potentially leading to inaccuracies. Researchers now show how to simultaneously determinethe parameters for visible light using periodic metal structures.Thomas Lepetit and Boubacar Kanté are in

Thomas Lepetit and Boubacar Kanté

Thomas Lepetit and Boubacar Kanté are in the Department of Electrical and Computer Engineering, University of California San Diego,

### The operation wavelength range is ~ 750-850 nm

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#### erc

![](_page_43_Picture_9.jpeg)

1 -1

### GSP-based phase-gradient metasurfaces: <u>skin-tight invisibility (carpet) cloak</u>!

#### **APPLIED OPTICS**

## An ultrathin invisibility skin cloak for visible light

![](_page_44_Figure_3.jpeg)

![](_page_45_Picture_0.jpeg)

## Can we implement optical stealth technologies, i.e., nullify backscattering (radar) cross section?

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![](_page_45_Picture_3.jpeg)

![](_page_46_Figure_0.jpeg)

#### Isotropic metasurfaces with uniform random phase response:

For 
$$r_{ms}(x', y') = a \exp[i\phi(x', y')]$$
 and PDF:  $p(\phi) = \frac{1}{2\pi}, -\pi < \phi < \pi$   
 $\Rightarrow$  then PDF of the far-field intensity:  $p(I) = \frac{1}{\overline{I}} \exp(-I/\overline{I})$ 

#### Similar to scattering from atmosphere (density fluctuations)!

![](_page_46_Picture_5.jpeg)

## GSP-based gradient metasurfaces: Random-phase optical metasurfaces

![](_page_47_Figure_1.jpeg)

**Figure 2.** Numerical modelling of scattering from ideal random-phase metasurfaces. Probability density function of the far-field intensity from an array of  $200 \times 200$  unit cells for excitation by (**a**) *x*-polarised incident light when unit cells feature random phases described by equation (3); (**b**) *x*-polarised incident light when the random phases of the unit cells are represented by a discrete random variable that takes on values ( $0, \pi/2, \pi, 3\pi/2$ ),

![](_page_47_Picture_3.jpeg)

## GSP-based gradient metasurfaces: Random-phase optical metasurfaces

![](_page_48_Figure_1.jpeg)

#### **Optical stealth effect can be achieved with proper design!**

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![](_page_48_Picture_4.jpeg)

Sci. Rep. 6, 28448 (2016).

# Instead of out look

**The Question: Can one realize one-shot** spectropolarimetry (i.e., complete plane-wave characterization)?

![](_page_49_Picture_2.jpeg)

## GSP arrays for chiroptical spectroscopy

![](_page_50_Figure_1.jpeg)

#### But complete polarimetry is not really an option $\otimes$

![](_page_50_Figure_3.jpeg)

![](_page_51_Figure_1.jpeg)

# Careful adjustment of beam shape and position are essential!

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![](_page_51_Picture_4.jpeg)

![](_page_51_Picture_5.jpeg)

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![](_page_52_Picture_1.jpeg)

#### **STILL IN WORKS...**

![](_page_52_Figure_3.jpeg)

![](_page_53_Figure_1.jpeg)

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![](_page_53_Picture_3.jpeg)

[-1,0,0]

![](_page_53_Figure_5.jpeg)

[0,0,-1]

![](_page_53_Figure_7.jpeg)

![](_page_53_Figure_8.jpeg)

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#### $d = 64 \ \mu m, LP90deg$ $\lambda = 710, 730, 750, 780 \ 800, 830, 850, 880, 900, 930, 950, 980, 1000 \ nm$

![](_page_54_Picture_1.jpeg)

## $d = 64 \ \mu m, \ LP0 deg \\ \lambda = 710, 730, 750, 780 \ 800, 830, 850, 880, 900, 930, 950, 980, 1000 \ nm$

![](_page_55_Picture_1.jpeg)

![](_page_56_Figure_1.jpeg)

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# Acknowledgements

# Thank you for your attention!

### The Danish Council for Independent Research European Research Council

![](_page_57_Picture_4.jpeg)